

example of diagonalization of Bogoliubov Hamiltonian

homogeneous system at rest $\phi_0(z) = \phi_0$, $\phi_0 \in \mathbb{R}$

$$H_{\text{Bog}} = \frac{1}{2} (\Delta^\dagger, \Delta) \eta \mathcal{L} \begin{pmatrix} \Delta \\ \Delta^\dagger \end{pmatrix}$$

$$\mathcal{L} = \begin{pmatrix} -\frac{\hbar^2}{2m} \nabla^2 + \mu_0 & \mu_0 \\ -\mu_0 & +\frac{\hbar^2}{2m} \nabla^2 - \mu_0 \end{pmatrix}$$

translational invariance \rightarrow eigenvectors $\begin{pmatrix} U \\ V \end{pmatrix} e^{ikz}$

$$\vec{\psi}_{+,k} = \begin{pmatrix} U_k \\ V_k \end{pmatrix} e^{ikz}, \quad E_k = \sqrt{\frac{\hbar^2 k^2}{2m} \left(\frac{\hbar^2 k^2}{2m} + 2\mu_0 \right)}$$

$$U_k \pm V_k = \left(\frac{\hbar^2 k^2 / 2m}{\frac{\hbar^2 k^2}{2m} \left(\frac{\hbar^2 k^2}{2m} + 2\mu_0 \right)} \right)^{\pm 1/2}$$

for each k

$$\vec{\psi}_{-,k} = \begin{pmatrix} V_k \\ U_k \end{pmatrix} e^{ikz}, \quad E_k = -\sqrt{\dots}$$

same U_k, V_k

$$\int \vec{\psi}_{+,k}^\dagger \eta \vec{\psi}_{+,k'} = (2\pi)^d \delta(k-k')$$

$$\int \vec{\psi}_{+,k}^\dagger \eta \vec{\psi}_{+,k'} = 0$$

$$\int \vec{\psi}_{+,k}^\dagger \eta \vec{\psi}_{-,k'} = -(2\pi)^d \delta(k-k')$$

decompose $\begin{pmatrix} \Lambda \\ \Lambda^\dagger \end{pmatrix}$ over eigenvectors of \mathcal{L} :

$$\hat{b}_n = \vec{v}_{+n}^\dagger \eta \begin{pmatrix} \Lambda \\ \Lambda^\dagger \end{pmatrix}$$

$$\hat{c}_n = \vec{v}_{-n}^\dagger \eta \begin{pmatrix} \Lambda \\ \Lambda^\dagger \end{pmatrix} = -\hat{b}_{-n}^\dagger$$

$$\begin{aligned} \begin{pmatrix} \Lambda \\ \Lambda^\dagger \end{pmatrix} &= \int \frac{dk}{(2\pi)^d} \vec{v}_{+n} \cdot \hat{b}_n - \vec{v}_{-n} \cdot \hat{c}_n = \\ &= \int \frac{dk}{(2\pi)^d} \begin{pmatrix} U_n \\ V_n \end{pmatrix} e^{ikz} \hat{b}_n + \begin{pmatrix} V_n \\ U_n \end{pmatrix} e^{-ikz} \hat{b}_n^\dagger \end{aligned}$$

$$\mathcal{L} \begin{pmatrix} \Lambda \\ \Lambda^\dagger \end{pmatrix} = \int \frac{dk}{(2\pi)^d} \omega_n^{\text{Bog}} \begin{pmatrix} U_n \\ V_n \end{pmatrix} e^{ikz} \hat{b}_n - \omega_n^{\text{Bog}} \begin{pmatrix} V_n \\ U_n \end{pmatrix} e^{-ikz} \hat{b}_n^\dagger$$

$$\frac{1}{2} (\Lambda^\dagger \Lambda) \eta \mathcal{L} \begin{pmatrix} \Lambda \\ \Lambda^\dagger \end{pmatrix} = \frac{1}{2} \int \frac{dk}{(2\pi)^d} \omega_n^{\text{Bog}} (\hat{b}_n^\dagger \hat{b}_n + \hat{b}_n \hat{b}_n^\dagger) =$$

$$= \text{cte} + \int \frac{dk}{(2\pi)^d} \hat{b}_n^\dagger \hat{b}_n \omega_n^{\text{Bog}}$$

Physical meaning of squeezed form of Bogoliubov vacuum

$$\hat{\rho}(x) = \hat{\psi}^\dagger(x) \hat{\psi}(x) \equiv |\phi(x)|^2 \left(N - \int dy \Delta^\dagger(y) \Delta(y) \right) + \sqrt{N} \left(\phi_0^\dagger(x) \Delta^\dagger(x) + \phi_0(x) \Delta(x) \right) + \Delta^\dagger(x) \Delta(x)$$

(homogeneous system) $= \frac{N}{V} - \frac{1}{V} \int dy \Delta^\dagger(y) \Delta(y) + \Delta^\dagger(x) \Delta(x) + \sqrt{N} \left(\Delta^\dagger(x) + \Delta(x) \right)$

$$\hat{\rho}(k \neq 0) = \int \frac{dx}{\sqrt{V}} e^{-ikx} \hat{\rho}(x) = \int \frac{dx}{\sqrt{V}} e^{-ikx} \left(\Delta^\dagger(x) \Delta(x) + \sqrt{N} \left(\Delta^\dagger(x) + \Delta(x) \right) \right)$$

$$= \sqrt{m} \int \frac{dx}{\sqrt{V}} e^{-ikx} \sum_{k'} \left(\frac{U_{k'}}{\sqrt{V}} e^{ik'x} b_{k'} + \frac{V_{k'}}{\sqrt{V}} e^{-ik'x} b_{k'}^\dagger + h.c. \right) + \text{higher-order terms}$$

$$= \sqrt{m} \cdot [U_n b_n + V_n b_{-n}^\dagger + h.c.] = \sqrt{m} (U_n + V_n) (b_n + b_{-n}^\dagger)$$

Analogously: $\hat{j}(x) = -\frac{i\hbar}{2m} (\hat{\psi}^\dagger \nabla \hat{\psi} - \nabla \hat{\psi}^\dagger \cdot \hat{\psi})$

$$\Rightarrow \hat{j}(k) = \frac{\sqrt{m} \hbar k}{2m} (U_n - V_n) (b_n - b_{-n}^\dagger) = \frac{\hbar k}{2m} \sqrt{m} \cdot (U_n - V_n)$$

Under free evolution:

$$\frac{d}{dt} \rho(k) = \sqrt{m} (U_n + V_n) \cdot (-i\omega_n \hat{b}_n + i\omega_{-n} \hat{b}_{-n}^\dagger) = -i\omega_n \sqrt{m} (U_n + V_n) (b_n - b_{-n}^\dagger) =$$

and $\text{div } \hat{j} \rightarrow \hbar \cdot \hat{j}(k) = \frac{\hbar^2 k^2}{2m} \sqrt{m} (U_n - V_n) (b_n - b_{-n}^\dagger)$

$$\text{So: } \frac{d}{dt} \rho(k) + \text{div } \hat{j}(k) = -i\sqrt{m} \cdot \left[(U_n + V_n) \omega_n - \frac{\hbar k^2}{2m} (U_n - V_n) \right] = 0$$

$$\begin{aligned} \langle e^{\dagger}(n) e(n) \rangle &= m (\omega_n + v_n)^2 \langle (b_n^{\dagger} + b_{-n}) (b_n + b_{-n}^{\dagger}) \rangle \\ &= m (\omega_n + v_n)^2 \left(\frac{2}{e^{\beta \hbar \omega_n} - 1} + 1 \right) = m \frac{\hbar^2 \omega_n^2 / 2m}{\omega_n} \left(\frac{2}{e^{\beta \hbar \omega_n} - 1} + 1 \right) \end{aligned}$$

$$\langle f^{\dagger}(n) f(n) \rangle = m \frac{\hbar^2 \omega_n^2}{4m^2} \frac{\omega_n}{\hbar^2 \omega_n^2 / 2m} \left(\frac{2}{e^{\beta \hbar \omega_n} - 1} + 1 \right) =$$

$$= \frac{m \omega_n}{2m} \left(\frac{2}{e^{\beta \hbar \omega_n} - 1} + 1 \right)$$

NOTE:

$$\begin{aligned} \langle e(n) e(-n) \rangle &= m (\omega_n + v_n)^2 \\ &\cdot \langle (b_n^{\dagger} + b_{-n}) (b_{-n}^{\dagger} + b_n) \rangle \\ &= \langle e^{\dagger}(n) e(n) \rangle \end{aligned}$$

for low n:

$$\langle e^{\dagger}(n) e(n) \rangle \rightarrow \frac{m \hbar \omega_n}{\omega_n^2} \frac{\hbar^2 \omega_n^2}{m}$$

$$\langle f^{\dagger}(n) f(n) \rangle \rightarrow \frac{m \hbar \omega_n}{m}$$

* At a given point x

$$\langle e(x) \cdot e(x) \rangle - \langle e(x) \rangle^2 = \langle \sqrt{m}^2 (\Delta(x) + \Delta^{\dagger}(x))^2 \rangle =$$

$$= m \cdot \langle (\Delta(x) + \Delta^{\dagger}(x))^2 \rangle =$$

$$= \left\langle m \cdot \sum_{n_1} \frac{(\omega_{n_1} + v_{n_1})}{\sqrt{V}} \cdot (b_{n_1} + b_{-n_1}^{\dagger}) \cdot \left(\frac{\omega_{n_1} + v_{n_1}}{\sqrt{V}} \right) (b_{n_1}^{\dagger} + b_{-n_1}) \right\rangle =$$

$$= \frac{m}{V} \cdot \sum_n (\omega_n + v_n)^2 \langle (b_n + b_{-n}^{\dagger})^2 \rangle = \frac{m}{V} \sum_n (\omega_n + v_n)^2 \left(\frac{2}{e^{\beta \hbar \omega_n} - 1} + 1 \right)$$

$$\langle f(x) f(x) \rangle = \frac{m}{V} \sum_n \left(\frac{\hbar \omega_n}{2m} \right)^2 (\omega_n - v_n)^2 \langle (b_n - b_{-n}^{\dagger})^2 \rangle =$$

$$= \frac{m}{V} \sum_n \left(\frac{\hbar \omega_n}{2m} \right)^2 (\omega_n - v_n)^2 \langle (b_n - b_{-n}^{\dagger})^2 \rangle$$

single particle $\rightarrow \langle j(x) j(x) \rangle$ has a factor $\left(\frac{\hbar}{2m}\right)^2$ more than
 $\langle \delta p(x) \delta p(x) \rangle$

interactions \rightarrow $\left\{ \begin{array}{l} \text{enhance } \langle j(x) j(x) \rangle \text{ by } (U_n - V_n)^2 \\ \text{suppress } \langle \delta p(x) \delta p(x) \rangle \text{ by } (U_n + V_n)^2 \end{array} \right.$



SQUEEZING effect

NOTE: Heisenberg is maintained as $(U_n + V_n)(U_n - V_n) = 1$.

NOTE 2: $\frac{1}{\sqrt{2}} (b_n + b_{-n}^\dagger)$ and $\frac{-i}{\sqrt{2}} (b_n - b_{-n}^\dagger)$
 \parallel \parallel
 X P

are "conjugate variables" in 2-mode squeezing

$X, P \rightarrow [X, P] = i$

see e.g. D.F. Walls and G.J. Milburn, "Quantum Optics",
 Springer, Berlin, 1994.