

(Not 25) quantum duchs - sketch of solution

1) generic wave equation in homogeneous geometry:

$$\mathcal{O}(k, \omega) h(k, \omega) = \mathcal{F}(k, \omega)$$

↳ satisfies $\mathcal{O}(k, \omega(k)) = 0$. or dispersion law
↳ applied force.

In this vicinity $\mathcal{O}(k, \omega) \approx \mathcal{O}'(k, \omega(k)) (\omega - \omega(k))$

$$\Rightarrow \mathcal{O}'(k, \omega(k)) (\omega - \omega(k)) h(k, \omega) = \mathcal{F}(k, \omega)$$

Back to time-domain $\omega \rightarrow +i \frac{\partial}{\partial t}$

$$(i \frac{\partial}{\partial t} - \omega(k)) h(k, t) = \underbrace{\frac{\mathcal{F}(k, t)}{\mathcal{O}'(k, \omega(k))}}_{\text{renormalised force } f(k, t)}$$

2) Uniform motion implies $f(r, t) = f(r - vt)$, so:

$$\begin{aligned} f(k, \omega) &= \int d^3r \int dt e^{-i\mathbf{k}\cdot\mathbf{r}} e^{i\omega t} f(r - vt) = \\ &= \int d^3r e^{-i\mathbf{k}\cdot(\mathbf{r}-v\mathbf{t})} f(r - vt) \underbrace{\int dt e^{-i\mathbf{k}\cdot v\mathbf{t}} e^{i\omega t}}_{2\pi \delta(\omega - \mathbf{k}\cdot\mathbf{v})} = \\ &= \underbrace{\int d^3r e^{-i\mathbf{k}\cdot\mathbf{r}} f(r)}_{f(k)} \end{aligned}$$

Massive particle (e.g. such) exchange $\Delta h, \Delta \omega$ into wave

$P \rightarrow P - \hbar \Delta k$ $E_n \rightarrow E_n - \hbar \Delta \omega$

$$\frac{P^2}{2m} \rightarrow \frac{(P - \hbar \Delta k)^2}{2m} = \frac{P^2}{2m} - \frac{P \cdot \hbar \Delta k}{m} + \frac{\hbar^2 \Delta k^2}{2m}$$

small.

$$\approx \frac{P^2}{2m} - \hbar \Delta k \cdot v = E_n - \hbar \Delta \omega$$

$\Rightarrow \Delta \omega = \Delta k \cdot v$ relate energy and momentum of emitted wave.

3+4) $\omega(k) - v \cdot k = 0 \Rightarrow \sqrt{gk} = k \cdot v = v k_x$

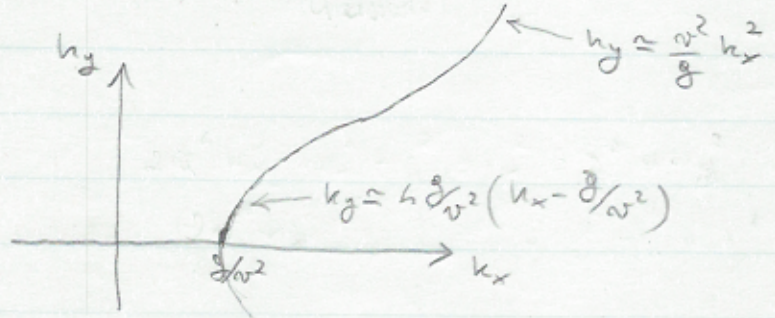
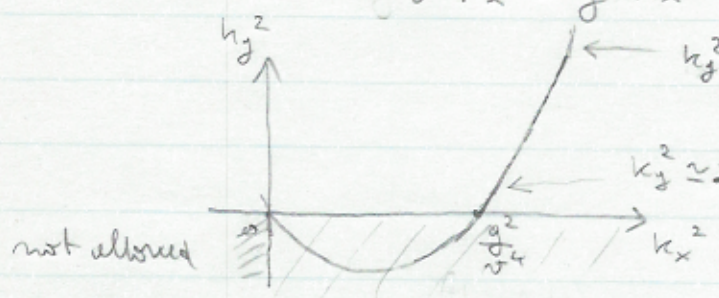
$$g \sqrt{k_x^2 + k_y^2} = v^2 k_x^2 \quad (\text{only for } k_x \geq 0)$$

$$v^4 k_x^4 - g^2 k_x^2 = g^2 k_y^2$$

$$k_y^2 \approx \frac{v^4}{g^2} k_x^4 \quad (\text{for } k_x^2 \gg \frac{g^2}{v^4})$$

$$k_y^2 \approx 2(k_x^2 - \frac{g^2}{v^4}) = 2(k_x + \frac{g}{v^2})(k_x - \frac{g}{v^2})$$

$$\approx 4 \frac{g}{v^2} (k_x - \frac{g}{v^2})$$

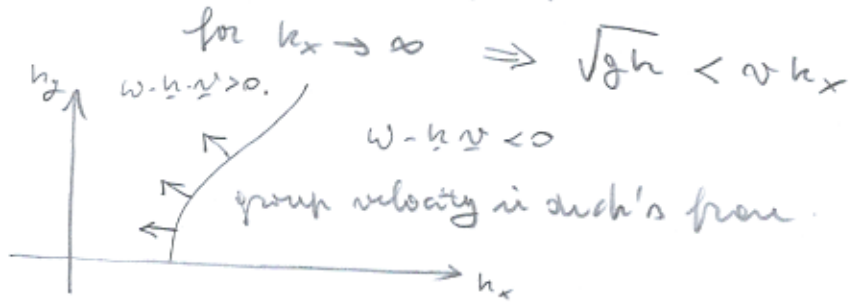


In lab frame $v_g = \nabla_n \omega$

In duck frame $v_g' = \nabla_n \omega - v = \nabla_n (\omega - kv)$,

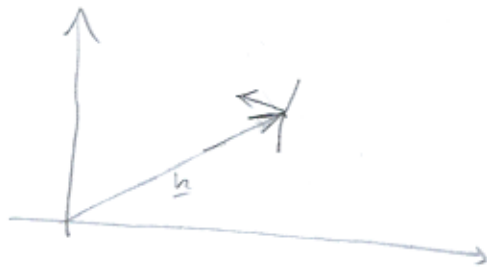
which is orthogonal to set $\omega - kv = 0$

To determine direction:



Roles of phase and group velocities explained in 1202.3494

In brief:



$k \in$ locus of $\omega - kv = 0 \rightarrow$ fringe pattern $\perp k$

\rightarrow extends in direction v_g' from duck



Angle of v_g' determined by $\frac{dk_y}{dk_x} = \frac{d}{dk_x} \frac{1}{g} \sqrt{v^4 k_x^4 - g^2 k_x^2} = \frac{2v^4 k_x^3 - g^2 k_x}{g \sqrt{v^4 k_x^4 - g^2 k_x^2}}$

$\min\left(\frac{dk_y}{dk_x}\right)$ is at $k_x^2 = \frac{3g^2}{2v^4}$

equal to $\sqrt{3}$.

So $v_g' \perp (1, \frac{dk_y}{dk_x}) \Rightarrow (v_g')_y = -\frac{1}{\frac{dk_y}{dk_x}} (v_g')_x$

$\max(\tan \theta_g) = \frac{v_g' y}{v_g' x} = -\frac{1}{\sqrt{3}}$, i.e. $\max \theta_g = 19.5^\circ$

In general:

for $w(u) \sim \ln|z|$

rescale $v \rightarrow \alpha v$

locus of $(w - kv) = 0$ is invariant

if $k \rightarrow \alpha^\beta k$ with $\beta z = \beta + 1$,
i.e. $\beta = \frac{1}{z-1}$

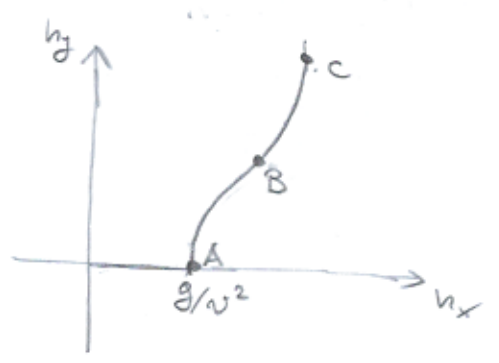
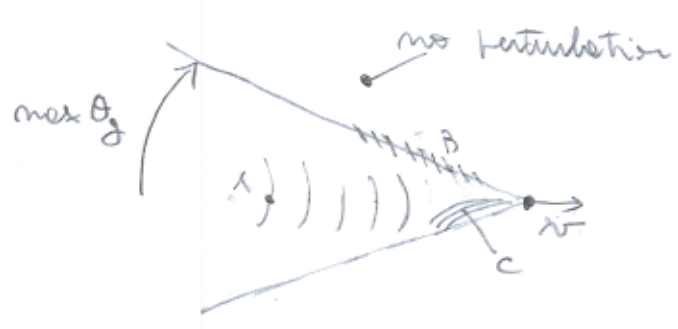
For our $z = 1/2$ case $\rightarrow \beta = -2$

(corresponds to threshold $k_{x,thr} = g/v^2$)

\Rightarrow rescaling of k does not change critical angle and fringe pattern, just spatial scale $r \rightarrow \alpha^\beta r$

For $z=1 \rightarrow$ no solution for $\beta \rightarrow$ critical angle varies with v

5) Shape of wake



From picture $\lambda_A \sim$ size arch $\sim 30 \text{ cm} = l_D$

$$\frac{4}{2\pi(g/v^2)} \Rightarrow \frac{2\pi v^2}{g} = l_D \Rightarrow v = \left(\frac{g l_D}{2\pi}\right)^{1/2} \approx 0,6 \text{ m/s}$$

(reasonable!)