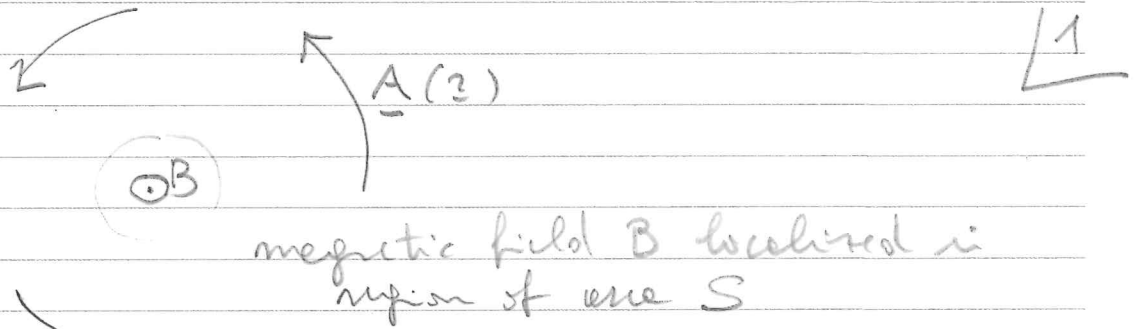


Aharonov - Bohm effect

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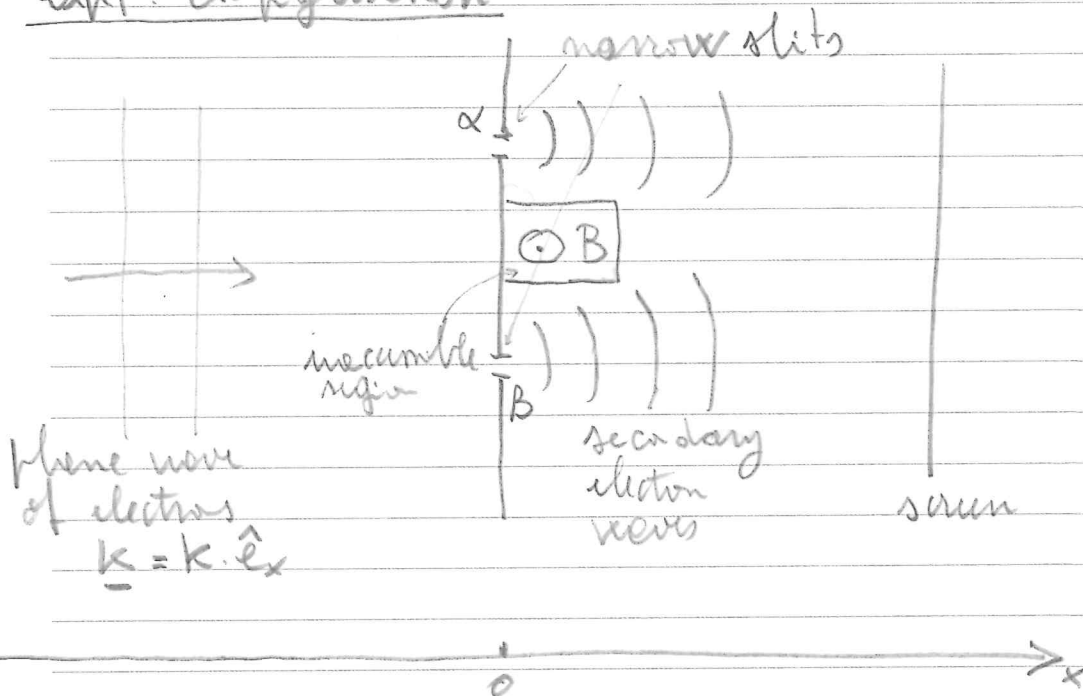
magnetic field B localized in region of area S

electrons can not penetrate this region

$$\underline{A}(z) = \frac{1}{2\pi r} \Phi \cdot \hat{e}_\theta, \quad \Phi = B \cdot S \text{ magnetic flux.}$$

$$\left[\text{from Stokes' thm: } \oint \underline{A}(z) \cdot d\underline{z} = \Phi = \int_S d^2\sigma \cdot \underline{B}(z) \right]$$

expt. configuration



gauge change

$$A' = A + \nabla\lambda \quad ; \quad \phi' = \phi - \frac{1}{c} \frac{\partial\lambda}{\partial t}$$

$$\psi' = \psi e^{iq\lambda/\hbar c}$$

$$\lambda = -\frac{1}{2\pi} \Phi \cdot \vartheta \rightarrow \nabla\lambda = -\frac{1}{2\pi} \Phi \frac{\hat{e}_\vartheta}{r}$$

$$A' = A + \nabla\lambda = \frac{1}{2\pi r} \Phi \hat{e}_\vartheta - \frac{\Phi}{2\pi r} \hat{e}_\vartheta = 0$$

but: $e^{iq\lambda/\hbar c} = \exp\left(-\frac{iq}{\hbar c} \frac{\Phi}{2\pi} \vartheta\right) =$

$$= \exp\left(-i \frac{q}{2\pi\hbar c} \Phi \vartheta\right)$$

is not continuous @ $\vartheta = 0$ vs. 2π

$$\frac{\Phi}{\Phi_0} = \frac{\Phi}{2\pi\hbar c/q} = \text{integer}$$

unit magnetic flux.

For $x \leq 0$: gauge change continuous as $\pi/2 < \vartheta < 3\pi/2$

For $x \geq 0$: discontinuity @ $\vartheta = 0, 2\pi$

$\text{define } \lambda' = -\frac{1}{2\pi} \Phi \vartheta \text{ for } -\frac{\pi}{2} > \vartheta > 0$	$\left. \begin{array}{l} \text{continuous} \\ \text{@ } \vartheta = 0, 2\pi \end{array} \right\}$
$= -\frac{1}{2\pi} \Phi (\vartheta - 2\pi) \text{ for } 2\pi > \vartheta > \frac{3\pi}{2}$	
	$\left. \begin{array}{l} \text{Discont.} \\ \text{moved} \\ \text{@ } \pi \end{array} \right\}$

for both gauge charges $A' = 0$

$$\psi_1 = e^{i q \vec{A} \cdot \vec{r} / \hbar c} \psi \quad \text{in } x \leq 0$$

in particular $\psi_1(\alpha) = \exp\left(-\frac{i q \Phi}{2\pi \hbar c} \frac{\pi}{2}\right) \psi(\alpha)$

$$\psi_1(\beta) = \exp\left(-\frac{i q \Phi}{2\pi \hbar c} \frac{3\pi}{2}\right) \psi(\beta)$$

at the slits located @ α, β

In $x \leq 0$ $\psi_1(z) = e^{i k x} \psi_0$

$$\psi_2(\alpha) = \exp\left(-\frac{i q \Phi}{2\pi \hbar c} \frac{\pi}{2}\right) \psi(\alpha) = \psi_1(\alpha)$$

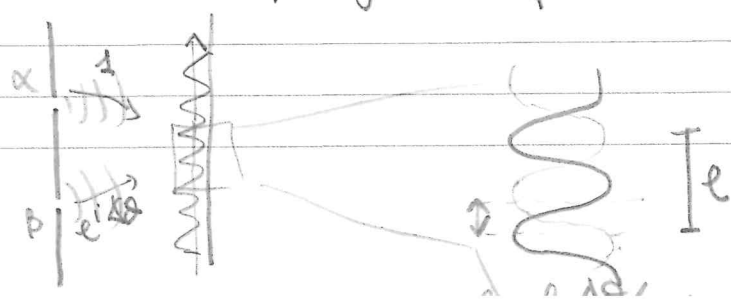
$$\psi_2(\beta) = \exp\left(-\frac{i q \Phi}{2\pi \hbar c} \left(\frac{\pi}{2} - 2\pi\right)\right) \psi(\beta) = e^{i q \Phi / \hbar c} \psi_1(\beta)$$

In $x \geq 0$: secondary waves have relative phase

$$\Delta\theta = \frac{q\Phi}{\hbar c} = 2\pi \frac{q\Phi}{2\pi \hbar c} = 2\pi \frac{\Phi}{\Phi_0}$$

As $A' = 0 \rightarrow$ standard two-slit interference

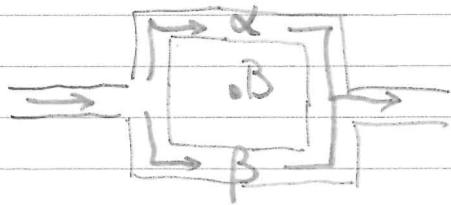
fringes shifted of $\Delta\theta = 2\pi \Phi / \Phi_0$



For superconductors :

elementary object "Cooper pair" \rightarrow charge $2q$

SQUID device



interference of current along paths α, β with

$$\text{relative phase } \Delta\theta = 2\pi \frac{\Phi}{2\pi\hbar c / 2q} = 4\pi \frac{\Phi}{\Phi_0}$$

Effective Josephson current :

$$J = J_\alpha \cdot \sin(\theta_\alpha) + J_\beta \sin(\theta_\beta)$$

$$\begin{aligned} J_\alpha &= J_\beta \\ &= J_0 (\sin \theta_\alpha + \sin \theta_\beta) = J_0 \left(\sin\left(\theta_\alpha + \frac{\Delta\theta}{2}\right) + \sin\left(\theta_\alpha - \frac{\Delta\theta}{2}\right) \right) \\ &= J_0 \cdot 2 \sin \theta_\alpha \cos \frac{\Delta\theta}{2} \end{aligned}$$

$$= J_0 \underbrace{2 \cos\left(\frac{\Delta\theta}{2}\right)} \cdot \sin \theta_\alpha$$

effective Josephson current through SQUID