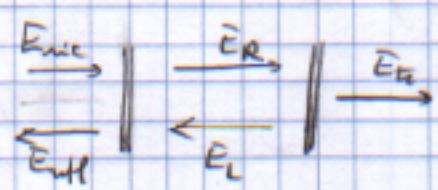


Reformulation in terms of field dynamics.



internal field :

$$E_R(\omega) = \frac{t}{1 - r^2 e^{2i\omega L/c}} E_{inc}(\omega)$$

$$E_L(\omega) = r e^{-i\omega L/c} E_R(\omega) = \frac{tr}{1 - r^2 e^{2i\omega L/c}} e^{-i\omega L/c} E_{inc}(\omega)$$

Around single resonance at  $\omega_0$ ,  $\Delta\omega = \omega - \omega_0$  :

$$E_R(\omega) \approx \frac{t}{1 - r^2 e^{2i\Delta\omega L/c}} E_{inc}(\omega) \approx \frac{t}{1 - r^2 - 2i\Delta\omega L/c r^2} E_{inc}(\omega) = \frac{i \frac{c}{2L}}{\omega - \omega_0 + i \frac{c}{2L} (1 - r^2)} E_{inc}(\omega)$$

i.e.:

$$(\omega - \omega_0 + i \frac{c}{2L}) E_R(\omega) = i \frac{c}{2L} t \cdot E_{inc}(\omega)$$

in real-time, this gives :

$$\left( i \frac{d}{dt} - \omega_0 - i \frac{c}{2L} \right) E_R(t) = i \frac{c}{2L} t E_{inc}(t)$$

$$\left( \frac{d}{dt} E_R = -i(\omega_0 - \frac{c}{2L}) E_R + \frac{c}{2L} t E_{inc}(t) \right)$$

$$E_L(t) = E_R(t - \frac{2L}{c}), \quad E_R(t) = t \cdot E_R(t - \frac{L}{c})$$

\* forced harmonic oscillator

- natural frequency  $\omega_0$
- damping  $\Gamma$
- forcing term proportional to incident field, coupling coefficient proportional to mirror  $t$ .

Model Hamiltonian (neglecting damping):

$\rightarrow$  non-Hamiltonian process,  
does not conserve phase-space volume.

$$\left\langle \hat{H} = \hbar \omega_0 \hat{a}^\dagger \hat{a} + i \left( \alpha t E_{inc}(t) \hat{a}^\dagger - \alpha^\dagger t^* E_{inc}^*(t) \hat{a} \right) \right\rangle$$

with external field operator  $\hat{E}_{L,R} = \mathcal{E}_{L,R} \hat{a}$ .

More complex processes can be included at the level of field dynamics and/or model Hamiltonian:

\*  $\chi^{(3)}$  optical nonlinearity:

$$\frac{d}{dt} E_R = \dots - i \omega_{nl} |E_R|^2 E_R$$

$$H^{(3)} = \hbar \omega_{nl} |\mathcal{E}_R|^2 \hat{a}^\dagger \hat{a}^\dagger \hat{a} \hat{a}$$

\* coupling with emitter:  $H_{int} = \hbar \Omega (a^\dagger b + a b^\dagger)$